

# Heating of planetesimals by SLRs $^{60}\text{Fe}$ & $^{26}\text{Al}$ and the effect on the water content of protoplanets

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## Context

Radiogenic sources of heating in the form of short-lived radioisotopes (SLRs) have been theorised since the 1950s to play an important role in planet formation [1]. SLRs were subsequently found to be the dominant source of planetesimal heating in the early solar system [2]. Radiogenic heating also results in loss of  $\text{H}_2\text{O}$  in a nascent planetesimal through vaporisation and out-gassing, which would significantly impact the liquid water content of the resultant protoplanet [3]. Whilst the SLR  $^{60}\text{Fe}$  did not provide sufficient heating in the early solar system, increased enrichment due to supernovae in star forming regions may result in a greater impact in other systems.

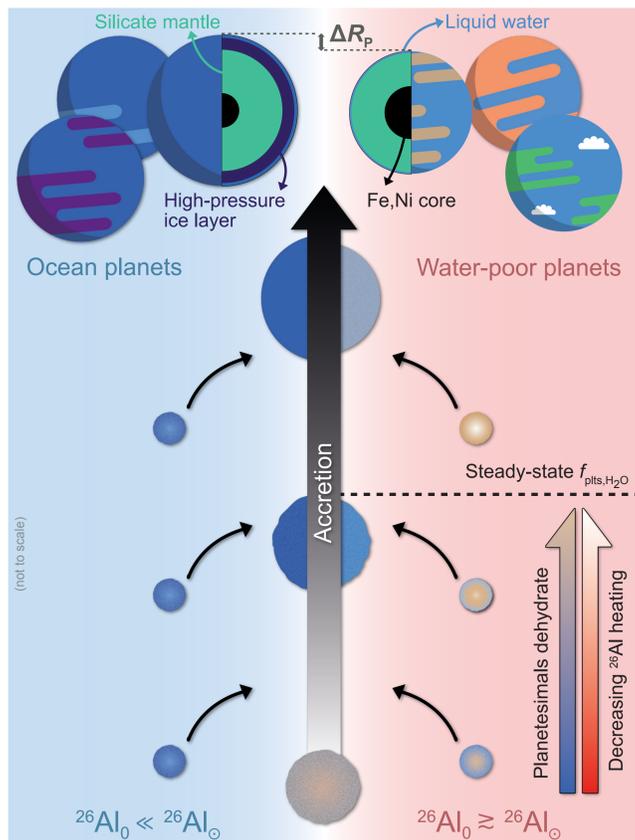


Figure 1. Planetary evolution due to desiccation [4].

## Methodology

The I2ELVIS [5] hydrodynamical code was used to perform multiple simulations to determine the influence SLR enrichment has on retained water in large planetesimals. Simulations had a maximum run time of 20 Myr, and planetesimals were assumed to form 1 Myr after CAI formation. The  $\text{H}_2\text{O}$  content within the planetesimal is tracked over time, and based on whether a cell's temperature has exceeded a threshold temperature of  $T_{\text{vap}} = 1223 \text{ K}$ . The final planetesimal  $\text{H}_2\text{O}$  abundance is then calculated from the fraction of cells that have exceeded  $T_{\text{vap}}$  ( $\mathcal{H}_f = N_{\text{wet}}/N_{\text{dry}}$ ). Radiogenic heating is simulated through introduction of a heating rate, given by the equation:

$$Q_{\text{SLR}}(t) = f_{\text{E,CI}} Z_{\text{SLR}} \frac{N_{\text{A}} E_{\text{SLR}} \lambda}{m_{\text{SLR}}} e^{-\lambda t}, \quad (1)$$

where  $f$  is the elemental abundance by mass,  $Z$  is the isotopic enrichment,  $N_{\text{A}}$  is Avogadro's constant,  $E$  is decay energy,  $\lambda$  is the decay constant,  $m$  is the atomic mass, and  $t$  is the elapsed time [6]. Enrichment is normalised to early solar system levels:

$$\Lambda_{\text{SLR}} = \frac{Z_{\text{SLR},*}}{Z_{\text{SLR},\odot}}, \quad (2)$$

where  $Z_{\text{SLR},*}$  is the isotopic enrichment of the system and  $Z_{\text{SLR},\odot}$  is the early solar system enrichment.

## Results – Iron Core Model

For the first model, a core of iron is surrounded by a mantle of hydrous silicates. The radius of this core relative to the planetesimal is described with variable  $\Psi = r_{\text{c}}/r_{\text{pl}}$ . Whilst this is a simple model to implement and simulate, this implies that core stratification occurs extremely quickly ( $\lesssim 1 \text{ Myr}$  from CAI formation) [7], however in the case of significant SLR enrichment this process could be accelerated. Due to this, this model as well as an undifferentiated model are simulated.

Core model,  $\Psi = 0.5$       Grain model,  $\Phi = 0.125$

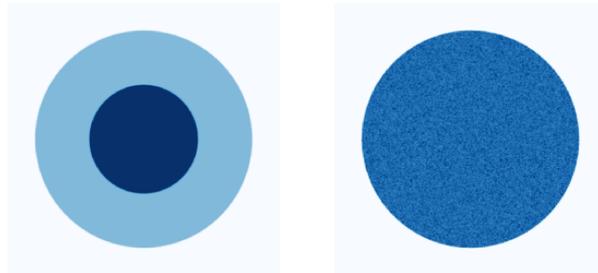


Figure 2. Comparison of core and grain composition models for similar iron volume fractions.

The parameter space exploration was of the isotopic enrichment of both  $^{26}\text{Al}$  and  $^{60}\text{Fe}$  for a planetesimal of radius 100 km.  $^{26}\text{Al}$  enrichment was varied from  $0 \leq \Lambda_{26\text{Al}} \leq 10$  and  $^{60}\text{Fe}$  enrichment was varied over  $0 \leq \Lambda_{60\text{Fe}} \leq 10^3$ .

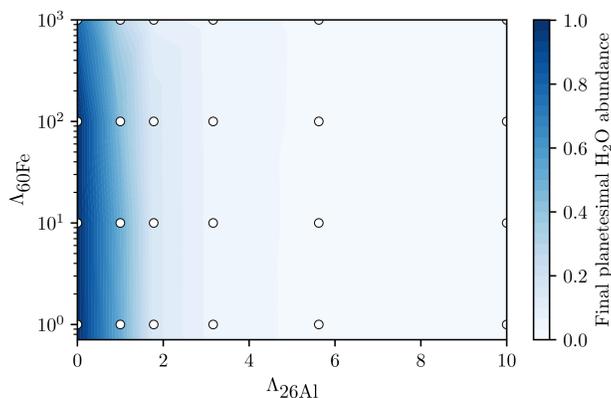


Figure 3. Remaining water fraction for simulations where  $\Psi = 0.25$ .

For simulations where  $\Psi = 0.25$  it was found that desiccation barely occurs even at an  $^{60}\text{Fe}$  enrichment of  $\Lambda_{60\text{Fe}} = 10^3 \times$ , where a final water fraction of  $\mathcal{H}_f = 0.796$  was found.

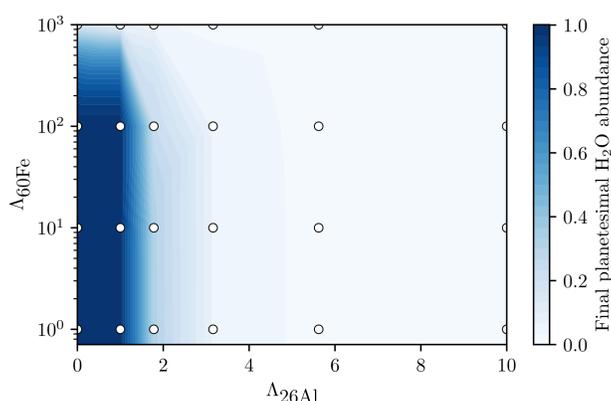


Figure 4. Remaining water fraction for simulations where  $\Psi = 0.5$ .

In simulations where  $\Psi = 0.5$ ,  $^{60}\text{Fe}$  enrichment is more influential on desiccation, but final water fraction in systems with high  $^{60}\text{Fe}$  enrichment and  $^{26}\text{Al}$  depletion are comparable to simulations where  $^{26}\text{Al}$  is enriched only slightly above early solar system estimates.

## Results – Iron Grain Model

The second model consists of a hydrous silicate body where iron marker “grains” are randomly interspersed. The amount of  $^{60}\text{Fe}$  grains is varied using a fraction,  $\Phi$  ( $N_{\text{Fe}}/N_{\text{T}}$ ). The main parameters explored in this set of simulations were  $\Phi$ ,  $^{60}\text{Fe}$  &  $^{26}\text{Al}$  enrichment.  $\Lambda_{60\text{Fe}}$  was varied from 0 to  $10^4$ , while  $\Phi$  was varied between 0 and 0.9. Subsets were performed where  $\Lambda_{26\text{Al}} = 0, 1 \& 10$ .

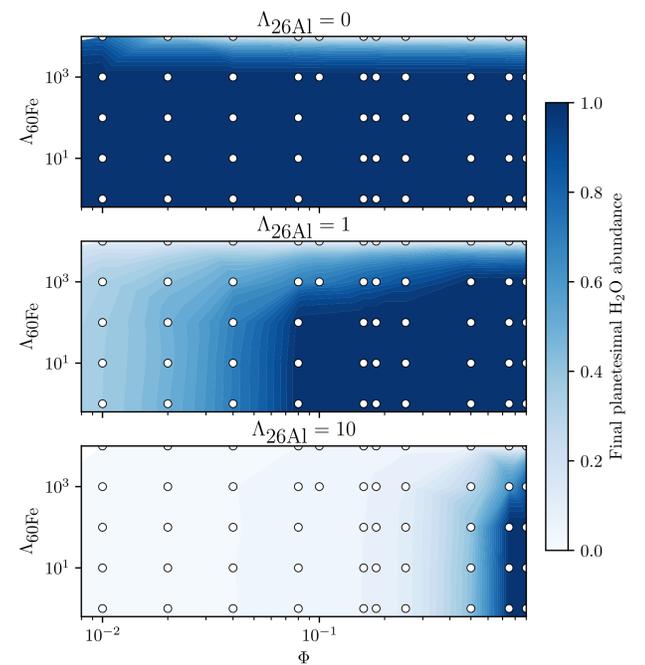


Figure 5. Remaining water fraction for simulations using the iron grain model.

Whilst a more accurate model overall, the grain model displays similar characteristics to the previous model.  $^{60}\text{Fe}$  is still not influential except in cases of extremely high enrichment ( $\Lambda_{60\text{Fe}} > 10^3$ ), high  $\Phi$  and total  $^{26}\text{Al}$  depletion. Even then, desiccation values are still far below what is observed with the core model. In cases of high  $^{26}\text{Al}$  enrichment, large quantities of iron grains can actively *inhibit* desiccation, as there is less  $^{26}\text{Al}$  to heat the planetesimal.

## Discussion

From these simulations it can be determined that  $^{60}\text{Fe}$  is a markedly less effective radiogenic heating source than  $^{26}\text{Al}$  during the planetesimal formation stage. Whilst some desiccation can occur as a result of  $^{60}\text{Fe}$  heating, it is significantly less drastic compared to even a slight enrichment of  $^{26}\text{Al}$  over early solar system estimates. Additionally, for extreme  $^{60}\text{Fe}$  enrichment the system may need to undergo extremely close or multiple supernovae encounters – which could inhibit the formation of the protoplanetary disk.

## References

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